

# On the exponential stability for a class of stochastic differential delay equations with fractional Brownian noises

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## Abstract

We investigate the asymptotic stability of the trivial solution of a stochastic differential delay equation which is driven by a fractional Brownian noise. Under the strong dissipative condition, we prove the exponential stability of the trivial solution under a small intensity of noise.

**Keywords:** fractional Brownian motion, stochastic differential equations (SDE), Young integral, exponential stability, random attractor.

## 1 Introduction

In this paper, we study the asymptotic stability of the following stochastic differential delay equation

$$dy(t) = [Ay(t) + f(y(t-r))]dt + Cy(t)dB^H(t), \quad (1.1)$$

given initial function  $\eta \in \mathcal{C}([-r, 0], \mathbb{R}^d)$ , where  $A, C \in \mathbb{R}^{d \times d}$  are matrices,  $r$  is a constant delay;  $f(0) = 0$  and  $f$  is a Lipschitz function, i.e.

$$\|f(y_1) - f(y_2)\| \leq C_f \|y_1 - y_2\|. \quad (1.2)$$

$B^H$  is an one-dimensional fractional Brownian motion [12] on a probability space  $(\Omega, \mathcal{F}, \mathbb{P})$  with the Hurst index  $H > 1/2$ . Equation (1.1) is understood and solved in the pathwise sense as a Young differential equation with delay

$$dy(t) = [Ay(t) + f(y(t-r))]dt + Cy(t)d\omega(t). \quad (1.3)$$

When  $r = 0$ , system (1.1) has the non-delay form, and the exponential stability of the trivial solution has been studied in recent papers [4], [7], [6] for Young equations. The stability problem for the Young differential delay equation (1.1) is however still in a fancy stage and is only studied in a few references, e.g. [11], although solutions of the Young differential equations have been well studied, [3], [5], or [6]. Since it is not sure how to apply the semigroup technique developed in [6] for the delay situation, our aim in this paper is to apply the techniques developed in [7] for nondelay equations to prove the exponential stability for the trivial solution of (1.1).

We present here a brief introduction to Young differential equations for Hölder continuous functions

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to match the delay situation (a detailed explanation can be found in [3] or [5]). For each  $0 < \alpha < 1$ , denote by  $C^\alpha([a, b], \mathbb{R}^d)$  the space of Hölder continuous functions on  $[a, b]$  equipped with the norm

$$\|u\|_{\infty, \alpha, [a, b]} := \|u\|_{\infty, [a, b]} + \|u\|_{\alpha, [a, b]}$$

where  $\|\cdot\|_{\infty, [a, b]}$  is the sup norm of continuous functions on  $[a, b]$  and

$$\|u\|_{\alpha, [a, b]} := \sup_{a \leq s < t \leq b} \frac{\|u(t) - u(s)\|}{(t - s)^\alpha}.$$

We also introduce the space

$$C^{0, \alpha}([a, b], \mathbb{R}^d) := \left\{ u \in C^\alpha([a, b], \mathbb{R}^d) \mid \lim_{h \rightarrow 0} \sup_{\substack{a \leq s < t \leq b, \\ |t-s| < h}} \frac{\|u(t) - u(s)\|}{(t - s)^\alpha} = 0 \right\}.$$

Given  $y \in C^\beta([a, b], \mathbb{R}^d)$  and  $\omega \in C^\nu([a, b], \mathbb{R})$ ,  $\beta + \nu > 1$ , it is well known that the *Young integral*  $\int_a^b y(t) d\omega(t)$  exists (see [13, p. 264–265], [5, p. 316]). This integral satisfies additive property by the construction, and the so-called Young-Loeve estimate [8, Theorem 6.8, p. 116]

$$\left\| \int_s^t y(u) d\omega(u) - y(s)[\omega(t) - \omega(s)] \right\| \leq K |t - s|^{\beta + \nu} \|y\|_{\beta, [s, t]} \|\omega\|_{\nu, [s, t]}, \quad (1.4)$$

where

$$K := (1 - 2^{1-\theta})^{-1}, \quad \theta := \beta + \nu > 1. \quad (1.5)$$

A result given in [1] shows that, under some regularity conditions of  $f$ , i.e.  $f$  is Lipschitz continuous, then there exists a unique solution  $y(\cdot, \omega, \eta)$  of (1.1) on  $\mathcal{C}^\beta([-r, T], \mathbb{R}^d)$  for  $H > \beta > 1 - H$  almost surely. A similar result can be found in [5]. In fact, (1.1) can be solved by induction in each interval  $[kr, (k+1)r]$  as a Young differential equation. Namely, given  $y_0 = \eta \in \mathcal{C}^\beta([-r, 0], \mathbb{R}^d)$ , then (1.3) derives

$$dy(kr + t) = [Ay(kr + t) + f(y((k-1)r + t))]dt + Cy(kr + t)d\theta_{kr}\omega(t), \quad \forall t \in [0, r].$$

The solution of (1.3) can be written by induction in the explicit form as  $y_0(\cdot) = \eta(\cdot) \in \mathcal{C}^\beta$  and

$$y(kr + t) = \Phi(t, \theta_{kr}\omega)y(kr) + \int_0^t \Phi(t - s, \theta_{kr+s}\omega)f(y((k-1)r + s))ds, \quad t \in [0, r] \quad (1.6)$$

for all  $k \geq 1$ , where  $\Phi(t, x)$  is the matrix solution of the linear Young differential equation

$$dz(t) = Az(t)dt + Cz(t)d\omega(t) \quad (1.7)$$

and satisfies  $\Phi(0, \omega) = Id$ .

Due to [2, Remark 2.3]  $\Phi(\cdot, \theta_{kr}\omega)y(kr) \in C^\beta([0, r], \mathbb{R}^d)$ . Hence (1.6) implies that  $y(kr + \cdot) \in C^\beta([0, r], \mathbb{R}^d)$ . By induction, one can prove that  $y(\cdot) \in C^\beta([-r, T], \mathbb{R}^d)$  for all  $T \geq -r$ . As a result,  $y_t(\cdot)$  belongs to  $C^\beta([-r, 0], \mathbb{R}^d)$  for all  $t \geq 0$ .

## 2 Exponential stability

In this section we would like to study the exponential stability of the zero solution of (1.1). We will work with the canonical space  $(\Omega, \mathcal{F}, \mathbb{P}, \theta)$  of fBm as introduced in [3]. Namely, for fixed  $1/2 < \nu' < H$ , denote by  $C_0^{\nu'}(\mathbb{R}, \mathbb{R})$  the space of all paths  $\omega$  which belong to  $C^{\nu'}(I, \mathbb{R})$  for all closed interval  $I \subset \mathbb{R}$  and receive value 0 at time 0. Then  $(\Omega, \mathcal{F}, \mathbb{P}, \theta)$  is the induced space on  $C_0^{\nu'}(\mathbb{R}, \mathbb{R})$  of

$(C_0(\mathbb{R}, \mathbb{R}), \mathcal{B}, \mathbb{P}, \theta)$  where  $(C_0(\mathbb{R}, \mathbb{R}))$  is the space of all continuous paths  $\omega$  on  $\mathbb{R}$  such that  $\omega(0) = 0$ ,  $\mathcal{B}$  is the Borel  $-\sigma$  algebra generated by the compact open topology,  $\mathbb{P}$  is the Wiener measure on  $\mathcal{B}$  generated by a fractional Brownian motion and  $\theta$  is the Wiener shift define by  $\theta_t \omega(\cdot) = \omega(t+\cdot) - \omega(\cdot)$ . Note that the Wiener shift is ergodic due to [9]. It is proved in [3] that equation (1.1) then generates a so-called *random dynamical system*  $\varphi : \mathbb{R}_+ \times (\Omega, \mathcal{F}, \mathbb{P}, \theta) \times C^{0,\beta}([-r, 0], \mathbb{R}^d) \rightarrow C^{0,\beta}([-r, 0], \mathbb{R}^d)$  given by

$$\varphi(t, \omega, \eta)(s) = y(t + s, \omega, \eta), \quad \forall t \in \mathbb{R}_+, s \in [-r, 0], \omega \in \Omega, \eta \in C^{0,\beta}([-r, 0], \mathbb{R}^d).$$

From now on, we fix  $\nu \in (1/2, \nu')$  and  $\beta \in (1-\nu, \nu)$ . It is known that  $B^H(\omega)|_{[a,b]} \in C^{0,\nu}([a, b], \mathbb{R})$ ,  $\forall \omega \in \Omega$  and  $E \left\| \|B^H\|_{\nu, [a,b]}^m \right\|$  is bounded for all  $[a, b] \subset \mathbb{R}$  and  $m > 0$  (see for instance [11]).

To obtain the main results we need the following lemmas.

**Lemma 2.1** *Assume that for fixed  $T$ ,  $y$  satisfies*

$$\|y\|_{\beta, [a,b]} \leq M \left[ (b-a)^{1-\beta} + (b-a)^{\nu-\beta} \|\omega\|_{\nu, [a,b]} \right] \left[ L + (b-a)^\beta \|y\|_{\beta, [a,b]} \right] \quad (2.1)$$

for all  $0 \leq a < b \leq T$ , in which  $M, L$  are constants. Then

$$\|y\|_{\beta, [0,T]} \leq L[2 \vee (4M)]^{1+\frac{1}{\nu}} T^{-\beta} \left( 1 + T^{1+\frac{1}{\nu}} + T^{1+\nu} \|\omega\|_{\nu, [0,T]}^{1+\frac{1}{\nu}} \right), \quad (2.2)$$

in which  $a \vee b := \max\{a, b\}$ .

*Proof:* Fix  $\mu = \frac{1}{2M}$ , we construct a sequence  $\{t_i\}$  such that  $t_0 = 0$  and

$$t_{i+1} = \inf\{t \geq t_i \mid (t - t_i) + (t - t_i)^\nu \|\omega\|_{\nu, [t_i, t]} \geq \mu\}.$$

Since  $|\|\omega\|_{\nu, [t_i, t']} - \|\omega\|_{\nu, [t_i, t]}| \leq \|\omega\|_{\nu, [t, t']}$  for all  $t_i \leq t \leq t'$  (see [5, p. 316]) and  $\omega \in C^{0,\nu}(I, \mathbb{R})$  for all closed interval  $I$ ,  $|\|\omega\|_{\nu, [t_i, t']} - \|\omega\|_{\nu, [t_i, t]}| \leq \|\omega\|_{\nu, [t, t']} \rightarrow 0$  as  $t' \rightarrow t$  (see [8, Corollary 5.31, p. 96]). It follows that for given  $t_i$  the function  $f(t) = (t - t_i) + (t - t_i)^\nu \|\omega\|_{\nu, [t_i, t]}$  is continuous on  $[t_i, +\infty)$ . Moreover,  $f$  is strictly increasing,  $f(t_i) = 0$  and  $\lim_{t \rightarrow +\infty} f(t) = +\infty$ . Therefore, there exists  $t_{i+1}$  such that

$$(t_{i+1} - t_i) + (t_{i+1} - t_i)^\nu \|\omega\|_{\nu, [t_i, t_{i+1}]} = \mu. \quad (2.3)$$

Assign  $N := \sup\{n \mid t_n \leq T\} < \infty$ .

For  $i = 0, \dots, N-1$ , (2.3) leads to

$$\mu = (t_{i+1} - t_i)^\nu \left[ (t_{i+1} - t_i)^{1-\nu} + \|\omega\|_{\nu, [t_i, t_{i+1}]} \right] \leq (t_{i+1} - t_i)^\nu \left( T^{1-\nu} + \|\omega\|_{\nu, [0,T]} \right).$$

This implies

$$t_{i+1} - t_i \geq \left( \frac{\mu}{T^{1-\nu} + \|\omega\|_{\nu, [0,T]}} \right)^{\frac{1}{\nu}},$$

and

$$T \geq \sum_{i=0}^{N-1} (t_{i+1} - t_i) \geq N \left( \frac{\mu}{T^{1-\nu} + \|\omega\|_{\nu, [0,T]}} \right)^{\frac{1}{\nu}}.$$

It follows that

$$N \leq T \left( \frac{T^{1-\nu} + \|\omega\|_{\nu, [0,T]}}{\mu} \right)^{\frac{1}{\nu}}. \quad (2.4)$$

By the assumption, for  $0 \leq i \leq N-1$  we have

$$\|y\|_{\beta, [t_i, t_{i+1}]} \leq M \left[ (t_{i+1} - t_i)^{1-\beta} + (t_{i+1} - t_i)^{\nu-\beta} \|\omega\|_{\nu, [t_i, t_{i+1}]} \right] \left[ L + (t_{i+1} - t_i)^\beta \|y\|_{\beta, [t_i, t_{i+1}]} \right]$$

or

$$\|y\|_{\beta, [t_i, t_{i+1}]} \leq \frac{ML}{1 - M\mu} \left[ T^{1-\beta} + T^{\nu-\beta} \|\omega\|_{\nu, [0, T]} \right].$$

We also have

$$\|y\|_{\beta, [t_N, T]} \leq \frac{ML}{1 - M\mu} \left[ T^{1-\beta} + T^{\nu-\beta} \|\omega\|_{\nu, [0, T]} \right].$$

Because  $[0, T] \subset \cup_{i=0}^{N-1} [t_i, t_{i+1}] \cup [t_N, T]$  and Hölder norm is of the superadditivity, we obtain

$$\begin{aligned} \|y\|_{\beta, [0, T]} &\leq \sum_{i=0}^{N-1} \|y\|_{\beta, [t_i, t_{i+1}]} + \|y\|_{\beta, [t_N, T]} \\ &\leq (N+1) \frac{ML}{1 - M\mu} \left[ T^{1-\beta} + T^{\nu-\beta} \|\omega\|_{\nu, [a, b]} \right]. \end{aligned} \quad (2.5)$$

Combining this with (2.4) we get

$$\begin{aligned} \|y\|_{\beta, [0, T]} &\leq 2ML \left[ 1 + T \left( \frac{T^{1-\nu} + \|\omega\|_{\nu, [0, T]}}{\mu} \right)^{\frac{1}{\nu}} \right] \left[ T^{1-\beta} + T^{\nu-\beta} \|\omega\|_{\nu, [0, T]} \right] \\ &\leq L(2M)^{1+\frac{1}{\nu}} T^{-\beta} \left[ \frac{1}{(2M)^{\frac{1}{\nu}}} + \left( T + T^\nu \|\omega\|_{\nu, [0, T]} \right)^{\frac{1}{\nu}} \right] \left[ T + T^\nu \|\omega\|_{\nu, [0, T]} \right] \\ &\leq L[1 \vee (2M)]^{1+\frac{1}{\nu}} T^{-\beta} \left[ 1 + \left( T + T^\nu \|\omega\|_{\nu, [0, T]} \right)^{\frac{1}{\nu}} \right] \left[ T + T^\nu \|\omega\|_{\nu, [0, T]} \right] \\ &\leq L[1 \vee (2M)]^{1+\frac{1}{\nu}} T^{-\beta} \left[ T + T^\nu \|\omega\|_{\nu, [0, T]} + \left( T + T^\nu \|\omega\|_{\nu, [0, T]} \right)^{1+\frac{1}{\nu}} \right]. \end{aligned} \quad (2.6)$$

Using the following inequalities

$$x^p + p - 1 \geq px, \quad \forall x \geq 0, p \geq 1, \quad (2.7)$$

and

$$(a+b)^p \leq (a^p + b^p) \max\{1, 2^{p-1}\}, \quad \forall a, b, p > 0, \quad (2.8)$$

we have

$$\begin{aligned} \|y\|_{\beta, [0, T]} &\leq L[1 \vee (2M)]^{1+\frac{1}{\nu}} T^{-\beta} \left[ \frac{(T + T^\nu \|\omega\|_{\nu, [0, T]})^{1+\frac{1}{\nu}}}{1 + \frac{1}{\nu}} + \frac{\frac{1}{\nu}}{1 + \frac{1}{\nu}} + (T + T^\nu \|\omega\|_{\nu, [0, T]})^{1+\frac{1}{\nu}} \right] \\ &\leq L[1 \vee (2M)]^{1+\frac{1}{\nu}} T^{-\beta} \left[ 2(T + T^\nu \|\omega\|_{\nu, [0, T]})^{1+\frac{1}{\nu}} + 1 \right] \\ &\leq L[2 \vee (4M)]^{1+\frac{1}{\nu}} T^{-\beta} \left( 1 + T^{1+\frac{1}{\nu}} + T^{1+\nu} \|\omega\|_{\nu, [0, T]}^{1+\frac{1}{\nu}} \right). \end{aligned} \quad (2.9)$$

The proof is completed.  $\square$

The following lemma is a Hölder-version of [7, Corrolary 3.5], which facilitates the proof of the stability of the delay system.

**Lemma 2.2** *Assuming that in (1.7)  $A$  is negative definite, i.e. there exists constant  $h_0 > 0$  such that*

$$\langle y, Ay \rangle \leq -h_0 \|y\|^2, \quad \forall y \in \mathbb{R}^d. \quad (2.10)$$

Then for all  $0 \leq s < t \leq r$ , the following inequalities hold

$$(i) \quad \|\Phi(t, \omega)\| \leq \exp\{-h_0 t + \|C\|tQ(r, \omega) + 2KG\|C\|\}, \quad (2.11)$$

$$(ii) \quad \|\Phi(t-s, \theta_s \omega)\| \leq \exp\{-h_0(t-s) + \|C\|(t-s)Q(r, \omega) + 2KG\|C\|\}, \quad (2.12)$$

where

$$G := \left( \max\{2, 8\|A\|, 16K\|C\|\} \right)^{1+\frac{1}{\nu}} \quad (2.13)$$

and

$$Q(t, \omega) := 2KG \left[ \frac{\nu(1+2KG)}{2KG} \|\omega\|_{\nu, [0, t]}^{\frac{1}{\nu}} + t^{\nu+\frac{1}{\nu}} \|\omega\|_{\nu, [0, t]} + t^{2\nu} \|\omega\|_{\nu, [0, t]}^{2+\frac{1}{\nu}} \right]. \quad (2.14)$$

*Proof:* (i) In the proof of [7, Theorem 3.4], we obtain

$$\begin{aligned} \log \|\Phi(t, \omega)\| &\leq -h_0 t + \left| \int_0^t \langle y(s), Cy(s) \rangle d\omega(s) \right| \\ &\leq -h_0 t + \|C\|t^\nu \|\omega\|_{\nu, [0, t]} \left( 1 + 2Kt^\beta \|\omega\|_{\beta, [0, t]} \right). \end{aligned}$$

in which  $y(t) = \frac{\Phi(t, \omega)}{\|\Phi(t, \omega)\|}$  satisfies

$$\|y\|_{\beta, [s, t]} \leq (2\|A\| \vee 4K\|C\|) \left[ (t-s)^{1-\beta} + (t-s)^{\nu-\beta} \|\omega\|_{\nu, [s, t]} \right] \left[ 1 + (t-s)^\beta \|y\|_{\beta, [s, t]} \right], \quad \forall s < t.$$

Apply Lemma 2.1 for  $M = 2\|A\| \vee 4K\|C\|$  and  $L = 1$ , we get

$$\|y\|_{\beta, [0, t]} \leq Gt^{-\beta} \left( 1 + t^{1+\frac{1}{\nu}} + t^{1+\nu} \|\omega\|_{\nu, [0, t]}^{1+\frac{1}{\nu}} \right).$$

Therefore, for  $0 \leq t \leq r$

$$\begin{aligned} \log \|\Phi(t, \omega)\| &\leq -h_0 t + \|C\|t^\nu \|\omega\|_{\nu, [0, t]} \left[ 1 + 2KG \left( 1 + t^{1+\frac{1}{\nu}} + t^{1+\nu} \|\omega\|_{\nu, [0, t]}^{1+\frac{1}{\nu}} \right) \right] \\ &\leq -h_0 t + \|C\|t^\nu \|\omega\|_{\nu, [0, t]} \left[ 1 + 2KG + 2KGt^{1+\frac{1}{\nu}} + 2KGt^{1+\nu} \|\omega\|_{\nu, [0, t]}^{1+\frac{1}{\nu}} \right] \\ &\leq -h_0 t + (1+2KG)\|C\|t^\nu \|\omega\|_{\nu, [0, t]} \\ &\quad + 2KG\|C\|t^\nu \|\omega\|_{\nu, [0, t]} \left( t^{1+\frac{1}{\nu}} + t^{1+\nu} \|\omega\|_{\nu, [0, t]}^{1+\frac{1}{\nu}} \right) \\ &\leq -h_0 t + (1+2KG)\|C\| \frac{t \|\omega\|_{\nu, [0, t]}^{\frac{1}{\nu}} + \frac{1}{\nu} - 1}{\frac{1}{\nu}} \\ &\quad + 2KG\|C\|t^\nu \|\omega\|_{\nu, [0, t]} \left( t^{1+\frac{1}{\nu}} + t^{1+\nu} \|\omega\|_{\nu, [0, t]}^{1+\frac{1}{\nu}} \right) \\ &\leq -h_0 t + 2KG\|C\| \\ &\quad + 2KG\|C\|t \left[ \frac{\nu(1+2KG)}{2KG} \|\omega\|_{\nu, [0, t]}^{\frac{1}{\nu}} + t^{\nu+\frac{1}{\nu}} \|\omega\|_{\nu, [0, t]} + t^{2\nu} \|\omega\|_{\nu, [0, t]}^{2+\frac{1}{\nu}} \right]. \end{aligned}$$

This implies (2.11).

(ii) For  $0 \leq s < t \leq r$ ,

$$\|\Phi(t-s, \theta_s \omega)\| \leq \exp \left\{ -h_0(t-s) + \|C\|(t-s)Q(t-s, \theta_s \omega) + 2KG\|C\| \right\}.$$

On the other hand

$$\begin{aligned} \|\theta_s \omega\|_{\nu, [0, t-s]} &= \sup_{0 \leq v < u \leq t-s} \frac{\|\omega(u+s) - \omega(v+s)\|}{|(u+s) - (v+s)|^\nu} \\ &= \sup_{s \leq v < u \leq t} \frac{\|\omega(u) - \omega(v)\|}{|u-v|^\nu} \leq \|\omega\|_{\nu, [0, r]}. \end{aligned}$$

As a result  $Q(t-s, \theta_s \omega) \leq Q(r, \omega)$ , which proves (2.12).  $\square$

To state our main result, let us introduce a Gronwall-Bellman type estimation, see the proof in [10].

**Proposition 2.3** *Let  $t_0 \in \mathbb{R}, 0 < t_0 \leq \infty, c \geq 0$  and  $a : [t_0, T] \rightarrow \mathbb{R}_+$  be locally integrable. Assume  $r \geq 0$ , and  $\tau : [t_0, T) \rightarrow \mathbb{R}_+$  is a measurable function such that  $t_0 - r \leq t - \tau(t), t_0 \leq t < T$ . If  $x : [t_0 - r, T) \rightarrow \mathbb{R}_+$  is Borel measurable and locally bounded such that*

$$x(t) \leq c + \int_{t_0}^t a(u)x(u - \tau)du, \quad t_0 \leq t < T, \quad (2.15)$$

then

$$x(t) \leq K \exp \left\{ \int_{t_0}^t \gamma(s)ds \right\}, \quad t_0 \leq t < T, \quad (2.16)$$

where the function  $\gamma : [t_0 - r, T) \rightarrow \mathbb{R}_+$  is locally integrable, and satisfies the characteristic inequality

$$a(t) \exp \left\{ - \int_{t-\tau(t)}^t \gamma(s)ds \right\} \leq \gamma(t), \quad t_0 \leq t < T, \quad (2.17)$$

and

$$K := \max \left\{ c \exp \left\{ \int_{t_0-r}^{t_0} \gamma(u)du \right\}, \sup_{t_0-r \leq s \leq t_0} \exp \left\{ \int_s^{t_0} \gamma(u)du \right\} \right\}. \quad (2.18)$$

When  $a$  and  $\tau$  are constant functions, we get the following corollary

**Corollary 2.4** *Let  $t_0 \in \mathbb{R}, t_0 < T \leq \infty$ , and  $c, a, \tau \geq 0$ . If  $y : [t_0 - \tau, T) \rightarrow \mathbb{R}_+$  is Borel measurable and locally bounded such that*

$$y(t) \leq c + \int_{t_0}^t ay(u - \tau)du, \quad t_0 \leq t < T, \quad (2.19)$$

then

$$y(t) \leq Ke^{\gamma(t-t_0)}, \quad t_0 \leq t < T, \quad (2.20)$$

where the nonnegative number  $\gamma$  satisfies the inequality

$$a \leq \gamma e^{\gamma\tau}, \quad (2.21)$$

and

$$K := \max \left\{ ce^{\gamma\tau}, \sup_{t_0-r \leq s \leq t_0} y(s)e^{\gamma(t_0-s)} \right\}. \quad (2.22)$$

Our first main result can be formulated as follows.

**Theorem 2.5** *Assume  $A$  is negative definite, i.e. there exists a  $h_0 > 0$  such that*

$$\langle y, Ay \rangle \leq -h_0 \|y\|^2. \quad (2.23)$$

and  $C_f < \frac{h_0}{2} e^{-\frac{h_0 r}{2}}$ . Then there exists  $\varepsilon > 0$  such that if  $\|C\| < \varepsilon$  the trivial solution of system (1.1) is forward exponentially stable almost surely, i.e. for any solution  $y$  of (1.1) we have

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log \|y_t(\omega, y_0)\|_{\infty, \beta, [-r, 0]} < 0, \quad (2.24)$$

for almost surely all the realizations  $\omega$ .

*Proof:* First, we fix  $\omega, \eta$  and denote by  $y(t)$ ,  $t \in [-r, \infty)$  the solution of (1.3). Since the assumption of the function  $f$ , it follows that

$$\|f(y)\| \leq C_f \|y\|. \quad (2.25)$$

Due to (1.6), (2.11), (2.12) and (2.25), for all  $t \in [0, r]$

$$\begin{aligned} \|y(kr+t)\| &\leq \exp\{-h_0 t + \|C\|tQ(r, \theta_{kr}\omega) + 2KG\|C\|\} \|y(kr)\| \\ &\quad + C_f \int_0^t \exp\{-h_0(t-s) + \|C\|(t-s)Q(r, \theta_{kr}\omega) + 2KG\|C\|\} \|y(kr+s-r)\| ds. \end{aligned} \quad (2.26)$$

For each  $k \in \mathbb{N}$  is fixed. Put

$$z^k(t) = \|y(kr+t)\| \exp\{h_0 t - \|C\|tQ(r, \theta_{kr}\omega)\}, t \in [-r, r]. \quad (2.27)$$

We have

$$\begin{aligned} z^k(t) &\leq e^{2KG\|C\|} \|y(kr)\| + e^{2KG\|C\|} C_f \int_0^t \exp\{h_0 s - \|C\|sQ(r, \theta_{kr}\omega)\} \|y(kr+s-r)\| ds \\ &\leq e^{2KG\|C\|} \|y(kr)\| + C_f \exp\{h_0 r - \|C\|rQ(r, \theta_{kr}\omega) + 2KG\|C\|\} \int_0^t z^k(s-r) ds. \end{aligned} \quad (2.28)$$

By the assumption  $C_f < \frac{h_0}{2} e^{-\frac{h_0 r}{2}}$ , there exists  $\varepsilon_1 > 0$  such that  $C_f < \frac{h_0}{2} \exp\{-\frac{h_0 r}{2} - 2KG\varepsilon_1\}$ . Then for  $\|C\| < \varepsilon_1$

$$C_f \exp\{h_0 r + 2KG\|C\|\} < \frac{h_0}{2} \exp\left\{\frac{h_0 r}{2} + 2KG(\|C\| - \varepsilon_1)\right\} < \frac{h_0}{2} e^{\frac{h_0 r}{2}}.$$

One can choose  $\gamma \in (0, \frac{h_0}{2})$  not depending on  $k$  and  $\omega$  such that

$$C_f \exp\{h_0 r - \|C\|rQ(r, \theta_{kr}\omega) + 2KG\|C\|\} \leq C_f \exp\{h_0 r + 2KG\|C\|\} \leq \gamma e^{\gamma r}.$$

Combining this with Corollary 2.4 we have for each  $t \in [0, r]$

$$z^k(t) \leq M_k e^{\gamma t}, \quad \text{or} \quad \|y(kr+t)\| \leq M_k \exp\{-h_0 t + \gamma t + \|C\|tQ(r, \theta_{kr}\omega)\}, \quad (2.29)$$

where

$$M_k = \max \left\{ e^{\gamma r + 2KG\|C\|} \|y(kr)\|, \sup_{s \in [-r, 0]} \|y(kr+s)\| \exp\{h_0 s - \gamma s - \|C\|sQ(r, \theta_{kr}\omega)\} \right\}. \quad (2.30)$$

Next, we introduce the function

$$\rho^k(t) = M_k \exp\{-h_0 t + \gamma t + \|C\|tQ(r, \theta_{kr}\omega)\}$$

then  $\|y(kr+t)\| \leq \rho^k(t)$  for all  $t \in [0, r]$  and

$$\rho^k(r) = M_k \exp\{-h_0 r + \gamma r + \|C\|rQ(r, \theta_{kr}\omega)\}.$$

Next we will prove that

$$M_k \leq \rho^{k-1}(r) \exp\{\gamma r + \|C\|rQ(r, \theta_{kr}\omega) + 2KG\|C\|\}. \quad (2.31)$$

Firstly, by the definition of  $M_k$  and (2.29), for  $s \in [-r, 0]$ ,

$$\begin{aligned}
& \|y(kr + s)\| \exp\{h_0s - \gamma s - \|C\|sQ(r, \theta_{kr}\omega)\} \\
& \leq \|y((k-1)r + s + r)\| \exp\{h_0(s+r) - \gamma(s+r)\} \exp\{-h_0r + \gamma r + \|C\|rQ(r, \theta_{kr}\omega)\} \\
& \leq \|y((k-1)r + s')\| \exp\{h_0s' - \gamma s'\} \exp\{-h_0r + \gamma r + \|C\|rQ(r, \theta_{kr}\omega)\} \\
& \leq M_{k-1} \exp\{\|C\|rQ(r, \theta_{(k-1)r}\omega)\} \exp\{-h_0r + \gamma r + \|C\|rQ(r, \theta_{kr}\omega)\},
\end{aligned} \tag{2.32}$$

in which  $0 \leq s + r = s' \leq r$ .

On the other hand,

$$M_{k-1} \exp\{-h_0r + \gamma r + \|C\|rQ(r, \theta_{(k-1)r}\omega)\} = \rho^{k-1}(r)$$

then

$$\sup_{[-r,0]} \|y(ks + r)\| \exp\{h_0s - \gamma s - \|C\|sQ(r, \theta_{kr}\omega)\} \leq \rho^{k-1}(r) e^{\|C\|rQ(r, \theta_{kr}\omega)}. \tag{2.33}$$

Secondly,

$$e^{\gamma r + 2KG\|C\|} \|y(kr)\| = e^{\gamma r + 2KG\|C\|} \|y((k-1)r + r)\| \leq \rho^{k-1}(r) e^{\gamma r + 2KG\|C\|}.$$

Hence

$$M_k \leq \rho^{k-1}(r) \exp\{\gamma r + \|C\|rQ(r, \theta_{kr}\omega) + 2KG\|C\|\}. \tag{2.34}$$

This implies

$$\begin{aligned}
\rho^k(r) & = M_k \exp\{-h_0r + \gamma r + \|C\|rQ(r, \theta_{kr}\omega)\} \\
& \leq \rho^{k-1}(r) \exp\{-h_0r + 2\gamma r + 2\|C\|rQ(r, \theta_{kr}\omega) + 2KG\|C\|\} \\
& \quad \dots \\
& \leq \rho^0(r) \prod_{j=1}^k \exp\{-h_0r + 2\gamma r + 2\|C\|rQ(r, \theta_{jr}\omega) + 2KG\|C\|\} \\
& = \rho^0(r) \exp\left\{\left[-(h_0 - 2\gamma)r + 2\|C\|r \frac{1}{k} \sum_{j=1}^k Q(r, \theta_{jr}\omega) + 2KG\|C\|\right]k\right\}.
\end{aligned} \tag{2.35}$$

Applying Birkhoff ergodic theorem

$$\begin{aligned}
& \lim_{k \rightarrow \infty} \left\{ -(h_0 - 2\gamma)r + 2KG\|C\| + \frac{2\|C\|r}{k} \sum_{j=1}^k Q(r, \theta_{jr}\omega) \right\} \\
& = -(h_0 - 2\gamma)r + 2KG\|C\| + 2\|C\|rE[Q(r, \omega)] =: -\lambda
\end{aligned} \tag{2.36}$$

holds for almost all  $\omega$ . Choose

$$\epsilon = \min \left\{ \epsilon_1, \frac{(h_0 - 2\gamma)r}{2KG + 2rE[Q(r, \omega)]} \right\}.$$

Then  $\lambda > 0$  provided that  $\|C\| < \epsilon$ . Moreover,  $\rho^0(r) = M_0 \exp\{-h_0r + \gamma r + \|C\|rQ(r, \omega)\}$ , in which

$$\begin{aligned}
M_0 & = \max \left\{ e^{\gamma r + 2KG\|C\|} \|\eta(0)\|, \sup_{s \in [-r,0]} \|y(s)\| \exp\{h_0s - \gamma s - \|C\|sQ(r, \omega)\} \right\} \\
& \leq \|\eta\|_{\infty, [-r,0]} \exp\left\{\gamma r + \|C\|Q(r, \omega) + 2KG\|C\|\right\}.
\end{aligned}$$

Hence

$$\rho^0(r) \leq \|\eta\|_{\infty, [-r,0]} \exp\left\{2\gamma r + 2\|C\|rQ(r, \omega) + 2KG\|C\|\right\}.$$



Hence when  $k$  is large enough, (2.36) deduce

$$\left[ -(h_0 - 2\gamma)r + 2KG\|C\| + 2\|C\|r \frac{1}{k} \sum_{j=1}^k Q(r, \theta_{jr}\omega) \right] k < \frac{-\lambda k}{2}. \quad (2.37)$$

As a result

$$\rho^k(r) \leq \rho^0(r) e^{-\frac{\lambda k}{2}}.$$

On the other hand, it is obvious that, for any  $0 \leq t \leq r$

$$\|y(kr + t)\| \leq \rho^k(t) \leq e^{h_0 r} \rho^k(r) \leq e^{-\frac{\lambda k}{2} + h_0 r} \rho^0(r).$$

Therefore

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log \|y_t\|_{\infty, [-r, 0]} < -\lambda/2. \quad (2.38)$$

Moreover, for any  $0 \leq t_1 < t_2 \leq r$ ,

$$\begin{aligned} & \|y(kr + t_2, \omega, \eta) - y(kr + t_1, \omega, \eta)\| \\ & \leq \left\| \left( \Phi(t_2, \theta_{kr}\omega) - \Phi(t_1, \theta_{kr}\omega) \right) y(kr, \omega, \eta) \right\| \\ & \quad + \left\| \int_{t_1}^{t_2} \Phi(t_2 - s, \theta_{s+kr}\omega) f(y(kr + s - r, \omega, \eta)) ds \right\| \\ & \quad + \left\| \int_0^{t_1} \left( \Phi(t_2 - s, \theta_{s+kr}\omega) - \Phi(t_1 - s, \theta_{s+kr}\omega) \right) f(y(kr + s - r, \omega, \eta)) ds \right\| \\ & \leq \|\Phi(t_2, \theta_{kr}\omega) - \Phi(t_1, \theta_{kr}\omega)\| \|y(kr, \omega, \eta)\| \\ & \quad + C_f \int_{t_1}^{t_2} \|\Phi(t_2 - s, \theta_{s+kr}\omega)\| \|y(kr + s - r, \omega, \eta)\| ds \\ & \quad + C_f \int_0^{t_1} \|\Phi(t_2 - s, \theta_{s+kr}\omega) - \Phi(t_1 - s, \theta_{s+kr}\omega)\| \|y(kr + s - r, \omega, \eta)\| ds \\ & \leq (t_2 - t_1)^\beta \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\beta, [0, r]} \\ & \quad + (t_2 - t_1) C_f \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \\ & \quad + C_f \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \int_0^{t_1} \left\| \left( \Phi(t_2, \theta_{kr}\omega) - \Phi(t_1, \theta_{kr}\omega) \right) \Phi^{-1}(s, \theta_{kr}\omega) \right\| ds \\ & \leq (t_2 - t_1)^\beta \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\beta, [0, r]} \\ & \quad + (t_2 - t_1) C_f \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \\ & \quad + (t_2 - t_1)^\beta r C_f \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\beta, [0, r]}. \end{aligned} \quad (2.39)$$

We then obtain for  $k$  large enough

$$\begin{aligned} \left\| \|y_{(k+1)r}\| \right\|_{\beta, [-r, 0]} & \leq \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\beta, [0, r]} \\ & \quad + r^{1-\beta} C_f \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \\ & \quad + r C_f \|y(kr + \cdot, \omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{kr}\omega)\|_{\beta, [0, r]} \\ & \leq e^{-\frac{\lambda(k-1)}{2} + h_0 r} \rho^0(r) \left[ 1 + C_f (r^{1-\beta} + r) \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \right] \|\Phi(\cdot, \theta_{kr}\omega)\|_{\infty, \beta, [0, r]}. \end{aligned} \quad (2.40)$$

On the other hand, due to [2, Proposition 2.4],  $\Phi$  and  $\Psi := \Phi^{-1}$  satisfy

$$\begin{aligned} d\Phi(t, \omega) & = A\Phi(t, \omega)dt + C\Phi(t, \omega)d\omega(t) \\ d\Psi(t, \omega) & = -A^T\Psi(t, \omega)dt - C^T\Psi(t, \omega)d\omega(t). \end{aligned}$$

Now, apply [6, Theorem 2.1] we get

$$\begin{aligned}\|\Phi(\cdot, \omega)\|_{\infty, [0, r]} &\leq e^{2\|A\|r} \left(\frac{K+2}{K+1}\right)^{1+r[2(K+1)\|C\|\|\omega\|_{\nu, [0, r]}]^{1/\nu}} \\ \|\Phi^{-1}(\cdot, \omega)\|_{\infty, [0, r]} &\leq e^{2\|A\|r} \left(\frac{K+2}{K+1}\right)^{1+r[2(K+1)\|C\|\|\omega\|_{\nu, [0, r]}]^{1/\nu}}.\end{aligned}$$

Moreover for all  $s, t \in [0, r]$

$$\begin{aligned}\|\Phi(\cdot, \omega)\|_{\beta, [s, t]} &\leq (\|A\| \vee K\|C\|) \left[ (t-s)^{1-\beta} + (t-s)^{\nu-\beta} \|\omega\|_{\nu, [s, t]} \right] \times \\ &\quad \times \left[ \|\Phi(\cdot, \omega)\|_{\infty, [0, r]} + (t-s)^{\beta} \|\Phi(\cdot, \omega)\|_{\beta, [s, t]} \right] \\ &\leq (\|A\| \vee K\|C\|) \left[ (t-s)^{1-\beta} + (t-s)^{\nu-\beta} \|\omega\|_{\nu, [s, t]} \right] \left[ L + (t-s)^{\beta} \|\omega\|_{\beta, [s, t]} \right]\end{aligned}$$

with  $L = e^{2\|A\|r} \left(\frac{K+2}{K+1}\right)^{1+r[2(K+1)\|C\|\|\omega\|_{\nu, [0, r]}]^{1/\nu}}$ .

Due to Lemma 2.1

$$\|\Phi(\cdot, \omega)\|_{\beta, [0, r]} \leq L[2 \vee 4\|A\| \vee 4K\|C\|]^{1+\frac{1}{\nu}} r^{-\beta} \left( 1 + r^{1+\frac{1}{\nu}} + r^{1+\nu} \|\omega\|_{\nu, [0, r]}^{1+\frac{1}{\nu}} \right).$$

Since  $\limsup_{k \rightarrow \infty} \frac{\|\theta_{kr}\omega\|_{\nu, [0, r]}}{k} = 0$ , it is evident that

$$\limsup_{k \rightarrow \infty} \log \left( \left[ 1 + C_f(r^{1-\beta} + r) \|\Phi^{-1}(\cdot, \theta_{kr}\omega)\|_{\infty, [0, r]} \right] \|\Phi(\cdot, \theta_{kr}\omega)\|_{\infty, \beta, [0, r]} \right) = 0$$

for all  $\omega$  in a set  $\Omega'$  with probability 1. Combining this with (2.40),  $\|y_{(k+1)r}\|_{\beta, [-r, 0]} \leq e^{-\lambda k/4}$  for  $k$  large enough. Now for  $t \in [kr, (k+1)r]$

$$\|y_t\|_{\beta, [-r, 0]} \leq \|y_{kr}\|_{\beta, [-r, 0]} + \|y_{(k+1)r}\|_{\beta, [-r, 0]} \leq 2e^{-\lambda(k-1)/4}.$$

Therefore

$$\limsup_{t \rightarrow \infty} \frac{\log \|y_t\|_{\beta, [-r, 0]}}{t} < -\lambda/4 \quad (2.41)$$

From (2.38) and (2.41) we obtain (2.24). The proof is completed.  $\square$

In the second main result, we prove that the trivial solution of (1.1) is also exponentially stable almost surely but in the pullback sense.

**Theorem 2.6** *Under the assumptions of Theorem 2.5, there exists  $\varepsilon > 0$  such that if  $\|C\| < \varepsilon$  the trivial solution of system (1.1) is pullback exponentially stable almost surely with respect to the  $\|\cdot\|_{\infty, \beta, [-r, 0]}$  norm, i.e.*

$$\limsup_{t \rightarrow \infty} \frac{1}{t} \log \|y_t(\theta_{-t}\omega, y_0)\|_{\infty, \beta, [-r, 0]} < 0,$$

for almost surely all the realizations  $\omega$ .

*Proof:* Firstly, we write  $y(t, \omega, \eta)$  and  $\rho^k(t, \omega, \eta)$  to stress the dependence of  $y$ ,  $\rho^k$  in Theorem 2.5 on  $\omega$  and  $\eta$ . The proof of this theorem follows the arguments of Theorem 2.5 with some modifications. Namely, fix  $u \in [kr, (k+1)r]$  and replace  $\omega$  in (2.35) by  $\theta_{-u}\omega$ , we obtain

$$\begin{aligned}\rho^k(r, \theta_{-u}\omega, \eta) &\leq \rho^0(r, \theta_{-u}\omega, \eta) \exp \left\{ \left[ - (h_0 - 2\gamma)r + \frac{2\|C\|r}{k} \sum_{j=1}^k Q(r, \theta_{jr-u}\omega) + 2KG\|C\| \right] k \right\} \\ &\leq \rho^0(r, \theta_{-u}\omega, \eta) \exp \left\{ \left[ - (h_0 - 2\gamma)r + \frac{2\|C\|r}{k} \sum_{j=0}^{k-1} P(r, \theta_{-jr}\omega) + 2KG\|C\| \right] k \right\},\end{aligned}$$

where

$$P(r, \omega) := 2KG \left[ \frac{\nu(1+2KG)}{2KG} \|\omega\|_{\nu, [-r, r]}^{\frac{1}{\nu}} + r^{\nu+\frac{1}{\nu}} \|\omega\|_{\nu, [-r, r]} + r^{2\nu} \|\omega\|_{\nu, [-r, r]}^{2+\frac{1}{\nu}} \right].$$

Put

$$\epsilon' = \min \left\{ \epsilon_1, \frac{(h_0 - 2\gamma)r}{2KG + 2rE[P(r, \omega)]} \right\}$$

and consider  $\|C\| < \epsilon'$  then for almost all  $\omega$

$$\begin{aligned} & \lim_{k \rightarrow \infty} \left\{ -(h_0 - 2\gamma)r + 2KG\|C\| + \frac{2\|C\|r}{k} \sum_{j=0}^{k-1} P(r, \theta_{-jr}\omega) \right\} \\ & = -(h_0 - 2\gamma)r + 2KG\|C\| + 2\|C\|rE[P(r, \omega)] =: -\lambda' < 0. \end{aligned} \quad (2.42)$$

Proof similarly to Theorem 2.5, it is easy to see

$$\rho^0(r, \theta_{-u}\omega, \eta) \leq \|\eta\|_{\infty, [-r, 0]} \exp \left\{ 2\gamma r + 2\|C\|rP(r, \theta_{-kr}\omega) + 2KG\|C\| \right\}.$$

It follows that when  $k$  is large enough,

$$\left[ -(h_0 - 2\gamma)r + 2KG\|C\| + 2\|C\|r \frac{1}{k} \sum_{j=0}^{k-1} P(r, \theta_{-jr}\omega) \right] k < \frac{-\lambda'k}{2}, \quad (2.43)$$

and

$$\rho^0(r, \theta_{-u}\omega, \eta) \leq e^{\frac{\lambda'k}{4}}.$$

As a result

$$\rho^k(r, \theta_{-u}\omega, \eta) \leq \rho^0(r, \theta_{-u}\omega, \eta) e^{\frac{-\lambda'k}{2}} \leq e^{\frac{-\lambda'k}{4}}.$$

On the other hand, it is obvious that, for any  $0 \leq t \leq r$

$$\begin{aligned} \|y(kr + t, \theta_{-u}\omega, \eta)\| & \leq \rho^k(t, \theta_{-u}\omega, \eta) \\ & \leq e^{h_0r} \rho^k(r, \theta_{-u}\omega, \eta) \\ & \leq e^{-\frac{\lambda'k}{2} + h_0r}. \end{aligned}$$

Also, one can choose  $k$  is greater than  $n(\omega)$  large enough so that

$$\|y(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [(k-2)r, (k+1)r]} \leq e^{-\frac{\lambda'(k-2)}{2} + h_0r}. \quad (2.44)$$

Next, in the analog of (1.6), for  $u \in [kr, (k+1)r]$  and  $u' = u - r$

$$y(u' + t, \omega, \eta) = \Phi(t, \theta_{u'}\omega)y(u', \omega, \eta) + \int_0^t \Phi(t-s, \theta_{s+u'}\omega)f(y(u' + s - r, \omega, \eta))ds.$$

Therefore, for  $k \geq n(\omega)$  and  $u \in [kr, (k+1)r]$ , we have

$$y(u' + t, \theta_{-u}\omega, \eta) = \Phi(t, \theta_{-r}\omega)y(u', \theta_{-u}\omega, \eta) + \int_0^t \Phi(t-s, \theta_{s-r}\omega)f(y(u' + s - r, \theta_{-u}\omega, \eta))ds.$$

For any  $0 \leq t_1 < t_2 \leq r$ , by similar computations to (2.39), we obtain

$$\begin{aligned}
& \|y(u' + t_2, \theta_{-u}\omega, \eta) - y(u' + t_1, \theta_{-u}\omega, \eta)\| \\
\leq & \left\| \left( \Phi(t_2, \theta_{-r}\omega) - \Phi(t_1, \theta_{-r}\omega) \right) y(u', \theta_{-u}\omega, \eta) \right\| \\
& + \left\| \int_{t_1}^{t_2} \Phi(t_2 - s, \theta_{s-r}\omega) f(y(u' + s - r, \theta_{-u}\omega, \eta)) ds \right\| \\
& + \left\| \int_0^{t_1} \left( \Phi(t_2 - s, \theta_{s-r}\omega) - \Phi(t_1 - s, \theta_{s-r}\omega) \right) f(y(u' + s - r, \theta_{-u}\omega, \eta)) ds \right\| \\
\leq & \|\Phi(t_2, \theta_{-r}\omega) - \Phi(t_1, \theta_{-r}\omega)\| \|y(u', \theta_{-u}\omega, \eta)\| \\
& + C_f \int_{t_1}^{t_2} \|\Phi(t_2 - s, \theta_{s-r}\omega)\| \|y(u' + s - r, \theta_{-u}\omega, \eta)\| ds \\
& + C_f \int_0^{t_1} \|\Phi(t_2 - s, \theta_{s-r}\omega) - \Phi(t_1 - s, \theta_{s-r}\omega)\| \|y(u' + s - r, \theta_{-u}\omega, \eta)\| ds \\
\leq & (t_2 - t_1)^\beta \|\Phi(\cdot, \theta_{-r}\omega)\|_{\beta, [0, r]} \|y_u(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} \\
& + (t_2 - t_1) C_f \|y_{u-r}(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \\
& + C_f \|y_{u-r}(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} \int_0^{t_1} \left\| \left( \Phi(t_2, \theta_{-r}\omega) - \Phi(t_1, \theta_{-r}\omega) \right) \Phi^{-1}(s, \theta_{-r}\omega) \right\| ds \\
\leq & (t_2 - t_1)^\beta \|\Phi(\cdot, \theta_{-r}\omega)\|_{\beta, [0, r]} \|y_u(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} \\
& + (t_2 - t_1) C_f \|y_{u-r}(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \\
& + (t_2 - t_1)^\beta r C_f \|y_{u-r}(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} \|\Phi^{-1}(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{-r}\omega)\|_{\beta, [0, r]}.
\end{aligned}$$

Due to (2.44) it follows that

$$\begin{aligned}
\|y_u(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, \beta, [-r, 0]} &= \|y_u(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, [-r, 0]} + \|y_u(\cdot, \theta_{-u}\omega, \eta)\|_{\beta, [-r, 0]} \\
&\leq e^{-\frac{\lambda'(k-2)}{2} + h_0 r} \times \left[ 1 + \|\Phi(\cdot, \theta_{-r}\omega)\|_{\beta, [0, r]} \right. \\
&\quad + r^{1-\beta} C_f \|\Phi^{-1}(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \\
&\quad \left. + r C_f \|\Phi^{-1}(\cdot, \theta_{-r}\omega)\|_{\infty, [0, r]} \|\Phi(\cdot, \theta_{-r}\omega)\|_{\beta, [0, r]} \right].
\end{aligned}$$

Consequently,

$$\lim_{u \rightarrow +\infty} \frac{1}{u} \log \|y_u(\cdot, \theta_{-u}\omega, \eta)\|_{\infty, \beta, [-r, 0]} = -\frac{\lambda'}{2} < 0.$$

That means the trivial solution of system (1.1) is pullback exponentially stable almost surely.  $\square$

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